

TEORETICKÝ POPIS PLAZMATU

- ČÁSTICOVÝ

- nejednotlivé částice
- problém mnoha těles
- småžeh + makroskopické pole (ne) relativistickými

- KINETICKÝ POPIS

- modelovací funkce  $f(\vec{r}, \vec{v}, t)$
- Boltzmannova rovnice

- TRKUTINOVÝ MODEL

- více tekutin
- vodorovná tekutina MHD

POHYB NAB. ČÁSTIC VE VN. POLI

- nerelativistickými

Lorentzova síla  $m \frac{d\vec{v}}{dt} = q(\vec{E} + \vec{v} \times \vec{B})$

$-\vec{E} = \text{const}$ ,  $|\vec{B}| = 0$  - rovnoměrně zrychlený pohyb

$-\vec{E} = 0$ ;  $\vec{B} = \text{const}$

$m \frac{d\vec{v}}{dt} = q \vec{v} \times \vec{B}$  ; necht  $\vec{B} \equiv B \hat{z}$

↓ roby

$m \dot{v}_x = q v_y B$  ;  $m \dot{v}_y = -q v_x B$  ;  $m \dot{v}_z = 0$

$\frac{d}{dt} \dot{v}_x = \frac{qB}{m} \dot{v}_y = -\left(\frac{qB}{m}\right)^2 v_x$

$\frac{d}{dt} \dot{v}_y = -\frac{qB}{m} \dot{v}_x = -\left(\frac{qB}{m}\right)^2 v_y$

$\omega_c = \frac{|q|B}{m}$  cyklotronová frekvence (Larmorova, gyro-)

$v_x = v_{\perp} \cos(\pm \omega_c t + \phi) = v_{\perp} \cos(\omega_c t)$

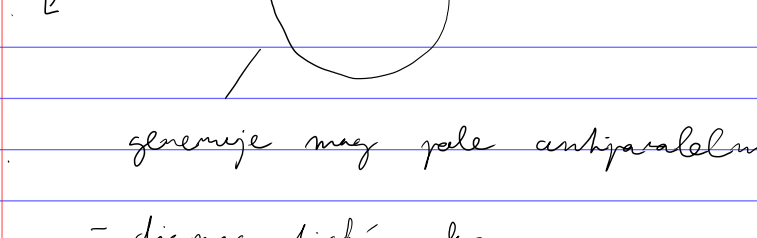
$v_y = \frac{\dot{v}_x}{\pm \omega_c} = \frac{1}{\pm \omega_c} v_{\perp} (-\sin(\omega_c t)) \omega_c = \mp v_{\perp} \sin(\omega_c t)$

$X = \int_0^t v_x dt + X_0 = \frac{v_{\perp}}{\omega_c} \sin(\omega_c t) + X_0$

$Y = \int_0^t v_y dt + Y_0 = \pm \frac{v_{\perp}}{\omega_c} \cos(\omega_c t) + Y_0$

$\frac{v_{\perp}}{\omega_c} = r_L$  Larmorova poloměr

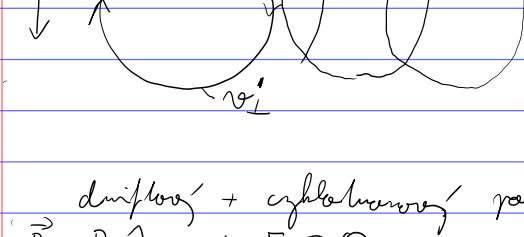
$r_L = \frac{m v_{\perp}}{q B}$



generuje mag pole antiparalelní  $\vec{B}$

- diamagnetické plazma
- pohyb pa rovnoběžnici  $\vec{v} = \vec{v}_{\parallel} + \vec{v}_{\perp}$

$\vec{E} = \text{const}$ ;  $\vec{B} = \text{const}$



driftový + cyklotronový pohyb

$\vec{B} = B \hat{z}$ ;  $E_y = 0$

$m \frac{d\vec{v}}{dt} = q(\vec{E} + \vec{v} \times \vec{B})$

$\dot{v}_z = \frac{q}{m} E_z$

$\dot{v}_x = \frac{q}{m} E_x + \omega_c v_y$

$\dot{v}_y = -\omega_c v_x$

$\frac{d}{dt}$

$\ddot{v}_x = -\omega_c^2 v_x$

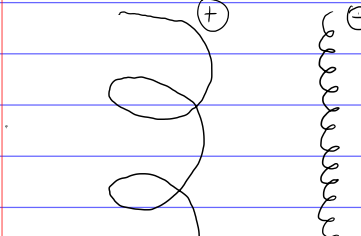
$\ddot{v}_y = +\omega_c \dot{v}_x = +\omega_c \left( \frac{q}{m} E_x + \omega_c v_y \right) =$

$= -\omega_c^2 \left( v_y + \frac{E_x}{B} \right)$

$\frac{d^2}{dt^2} \left( v_y + \frac{E_x}{B} \right) = -\omega_c^2 \left( v_y + \frac{E_x}{B} \right)$

$v_x = v_{\perp} \cos(\omega_c t)$

$v_y = \mp v_{\perp} \sin(\omega_c t) \rightarrow v_y = v_y' - \frac{E_x}{B} = \mp v_{\perp} \sin(\omega_c t) - \frac{E_x}{B}$



$\vec{v} = \vec{v}_{\perp} + \vec{v}_{\parallel}$

$m \frac{d}{dt} (\vec{v}_{\perp} + \vec{v}_{\parallel}) = q(\vec{E} + (\vec{v}_{\perp} + \vec{v}_{\parallel}) \times \vec{B})$

$\vec{E} + \vec{v}_{\parallel} \times \vec{B} = 0 \quad | \times \vec{B}$

$\vec{E} \times \vec{B} + (\vec{v}_{\parallel} \times \vec{B}) \times \vec{B} = 0 \quad | (a \times b) \times c = (c \cdot a)b - (c \cdot b)a$

$\vec{E} \times \vec{B} + (v_{\parallel} \cdot \vec{B}) \vec{B} - B^2 \vec{v}_{\parallel} = 0$

$\vec{v}_{\parallel} = \frac{\vec{E} \times \vec{B}}{B^2}$  drift ve zřícených polích E-cross-B drift

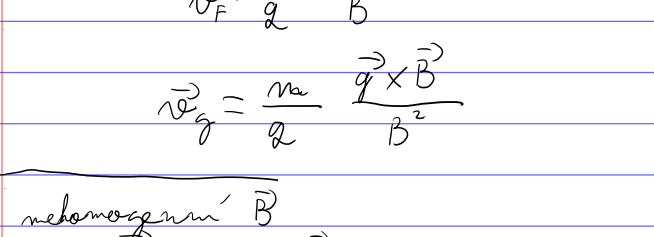
mezírání na q ani m

plát i pro obecnou sílu  $qE \rightarrow F$  (gravitaci, odstředivost)

$\vec{v}_F = \frac{1}{q} \frac{\vec{F} \times \vec{B}}{B^2}$

$\vec{v}_g = \frac{m}{q} \frac{\vec{g} \times \vec{B}}{B^2}$

nehomogenní  $\vec{B}$



$\nabla B \parallel \hat{y}$   $\vec{F} = q \vec{v} \times \vec{B}$

$\vec{F}_x = 0$  (symetrie)

$\vec{F}_y = -q v_x B_z$

$\vec{B} = \vec{B}_0 + (\vec{m} \cdot \nabla) \vec{B} + \dots$  plát pro  $m_L \ll \frac{B}{|\nabla B|}$

$B_z = B_0 + y \cdot \frac{\partial B_z}{\partial y} + \dots$

$\vec{F}_y = -q v_{\perp} \cos(\omega_c t) (B_0 \pm m_L \cos(\omega_c t) \frac{\partial B}{\partial y})$

$\overline{\cos(\omega_c t)} = 0 \quad \overline{\cos^2(\omega_c t)} = 1/2$

$\vec{F}_y = \mp q v_{\perp} m_L \frac{1}{2} \frac{\partial B}{\partial y}$

$v_{\perp} = \frac{1}{q} \frac{\vec{F} \times \vec{B}}{B^2} = \frac{1}{q} \frac{F_y}{|B|} \hat{x} = \mp \frac{v_{\perp} m_L}{B} \frac{1}{2} \frac{\partial B}{\partial y} \hat{x}$

$\vec{v}_{\nabla B} = \pm \frac{1}{2} v_{\perp} m_L \frac{\vec{B} \times \nabla B}{B^2}$  grad-B drift